PHILOSOPHICAL TRANSACTIONS.

I. On the Influence of Temperature on the Electric Conducting Power of Metals.

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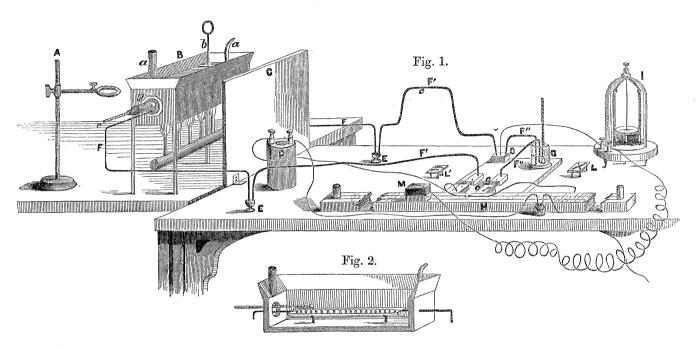
The results obtained by different observers in their researches on the influence of temperature on the electric conducting power of metals do not agree at all together. The differences in their results may be partly owing to their not having tested pure metals, and partly to their not having taken into consideration the fact that, when a wire of a pure metal is heated for the first time to 100° C., an alteration in the conducting power of the wire is observed on its again being cooled; in fact, it is necessary to keep the wire for several days at 100° before its conducting power, on again being cooled, becomes constant.

In the experiments we are about to detail we have taken great care to employ only pure metals, as well as a method and a disposition of the apparatus with which great accuracy could be obtained.

The method employed for the determination of the resistances is fully described in the 'Philosophical Magazine' for February 1857. Fig. 1 shows the disposition of the B is the trough in which the wires were heated: these were soldered to two thick copper wires F (4-5 millims, thick), bent as shown in the figure, and ending in the mercury-cups E, which were connected with the apparatus by two other copper wires, F', of the same thickness. C is a piece of board placed in such a manner as to prevent the heat of the trough from radiating on the apparatus. The mercurycups O are made of small blocks of wood, through which holes are bored just large enough to take the thick wires, and to the bottoms of which blocks amalgamated copper plates are fastened. Now it is clear that if the ends of the thick copper wires are filed flat, and well amalgamated, and the mercury-cups are filled with mercury, this method of connexion may be looked upon as a soldering of the copper plates to the wires, or, in other words, as a perfect connexion; for the wires may be removed as often as required, and on replacing them the same resistance is always observed. The wires F'', MDCCCLXII.

β**L**Δ11.

to which the normal wire (in the glass cylinder G) is soldered, are also 4-5 millims. thick. The reason why such thick wires were chosen was to make any difference in their resistance, caused by the change of temperature in the room or by the heating of the ends in the oil-bath, so small that no correction was necessary. This was proved to be the case by the following experiment:—After having soldered a wire in the trough to the ends of the thick copper wires, and determined its resistance with the normal wire generally used, the wire F' at e was heated with the 6-Bunsen burner much above



100° C., and the resistance of the circuit was again determined whilst the wire was at that temperature, when it was found to have increased only 0.08 per cent.; we did not, therefore, consider it necessary to make any correction for the increase of resistance caused by the heating of the ends of the thick wires in the trough. The resistance of the copper wires was determined at the ordinary temperature, and brought into calculation without further correction. Before the commencement of each series, all the ends of the wires dipping into the mercury-cups were carefully re-amalgamated. L, L' are the two commutators fitting into four mercury-cups at o.

The wire stretched on the board H is of german silver instead of copper, as was formerly described; its half-length was 4550 millims. The length of the board is about 1500 millims.; the wire, therefore, was wound backwards and forwards several times on the one side; this is not visible in the figure. By using normal wires of different resistances, and by choosing proper lengths of the wire to be tested, it was always possible to begin the observations with the block M within 100 millims. of the middle of the wire. Great care was taken to lift the block M off the wire when it was moved, in order to prevent as much as possible its wearing. It may be mentioned that, although we

generally worked with only one of the commutators, and therefore mostly used the one half of the wire, the zero-point of the wire only varied, during the whole of the experiments, which have taken almost a year to carry out, 3 millims. The zero-point was always determined before each series was begun. The distance the block M was moved when the resistance of a wire was determined, first at 0° and then at 100°, was, for pure metals in a solid state, about 800 millims., or about 8 millims. for 1°. As, however, the movement of the block M of 1 millim. caused a deflection of the needles of the galvanometer I of 20° to 30°, it is evident, with the apparatus employed, that the differences in the resistance of a wire to values less than those corresponding to 0°·1 C. can be accurately determined. Our results, moreover, prove this to be the case, as in many instances the difference between the observed and calculated conducting powers for the whole series do not amount to values equal to 0°·1 to 0°·2 C.

The trough B is a double one, the space between the inner and outer one being 20 millims. The dimensions of the inner trough were 400 millims long, 80 millims wide, and 80 millims deep. Through the ends of both two holes of about 20 millims wide were made, in which good corks were fitted, and through these passed the thick copper wires F; and also at one end a glass tube d, wide enough to allow the thermometer c to pass freely. A piece of india-rubber tubing, fitting over the glass tube d, and tightly round the thermometer, closed the tube, but allowed the thermometer to be moved either backward or forward with great ease. The tubes a are for filling the space between the inner and outer troughs with oil.

The wire to be tested lay in the trough, as shown in fig. 2, on a small glass tray, made by splitting a glass tube longitudinally, thereby preventing any possibility of its touching the trough, and also guarding it from being moved by the stirrer. A second trough, of somewhat smaller dimensions, was also used.

The use of an oil-bath for heating the wires has been objected to by a former observer*; it was therefore necessary to determine experimentally whether there was any real reason for the objection or not. He states that, as oil conducts electricity better on being heated than when cold, the differences between the conducting powers of cold and hot oil will materially affect the values obtained for the resistances of wire which had been determined at different temperatures in that liquid. In order to test the accuracy of this assertion, two copper plates of about 150 millims, diameter were connected, the one with the galvanometer, the other with a single Bunsen's cell; and to complete the circuit, this was connected with the galvanometer. A piece of filtering-paper, moistened with the olive-oil used, was placed between the copper plates, and these were pressed together with a weight. On completing the circuit not the slightest deflection of the needles was observed; the copper plates were then heated to above 100° C., and still no deflection was visible. To show that the connexions were good, a drop of water was put on the oiled paper; and immediately the needles of the galvanometer were sent with great violence to the stops. This proves that although oil may

^{*} Arndsten, Poggendorff's 'Annalen,' vol civ. p. 1.

have a higher conducting power when hot than cold, yet in either case it is so infinitely small, that it cannot influence the results obtained in the manner just described.

Again, it was proved in a former research* that the formula for the correction of conducting power for temperature of a wire, deduced from the observations made in an oil- or air-bath, were exactly the same. Thus the formula obtained for an annealed wire of the gold-silver alloy heated in the oil-bath was

$$\lambda = 15.052 - 0.01074t + 0.00000714t^2$$

and that for the same wire heated in an air-bath was

$$\lambda = 15.059 - 0.01077t + 0.00000722t^2$$
.

As, however, more accurate results may be obtained by experimenting in an oil-than in an air-bath, on account of the wires taking more readily the temperature of the bath, and of their being more rapidly cooled if heated by the current passing through them, we have chosen this manner of heating the wires in preference to the other.

As oil, and more especially oil when hot, attacks most wires to a degree which would render the observations valueless, we were obliged to varnish them. The best varnish for the purpose is a solution of shell-lac in alcohol. For instance, a hard-drawn copper wire, not varnished, loses in conducting power after having been heated in an oil-bath to 100°, but if varnished, increases. To show that varnishing has no effect on the results, we give in Table I. the conducting power of a hard-drawn gold wire, first not varnished, and then varnished. Each result is the mean of two observations.

	Not va	rnished.		Varnished.				
Т.	Conducting por		Difference.		Conducti	ng power.	73.100	
1.	Observed.	Calculated.	Dinerence.	Т.	Observed.	Calculated.	Difference.	
15·30 30·55 48·65 69·55 83·25 99·85 84·55 70·80 50·85 30·95 16·80	72.697 68.806 64.659 60.409 57.915 55.151 57.704 60.224 64.239 68.746 72.343	72·705 68·879 64·717 60·423 57·906 55·174 57·680 60·184 64·239 68·782 72·316	$\begin{array}{c} -0.008 \\ -0.073 \\ -0.058 \\ -0.014 \\ +0.009 \\ -0.023 \\ +0.024 \\ +0.040 \\ -0.036 \\ +0.027 \end{array}$	13.85 30.95 49.55 68.40 84.55 98.70 84.90 70.25 51.20 30.60 17.85	73·120 68·756 64·523 60·636 57·704 55·346 57·645 60·318 64·149 68·886 72·111	73·085 68·782 64·520 60·645 57·680 55·352 57·620 60·289 64·164 68·866 72·045	$\begin{array}{c} +0.035 \\ -0.026 \\ +0.003 \\ -0.009 \\ +0.024 \\ -0.006 \\ +0.025 \\ +0.029 \\ -0.015 \\ +0.020 \\ +0.066 \end{array}$	

TABLE I.

The formula deduced from the observations, and from which the conducting powers were calculated, was

$$\lambda = 76.838 - 0.27973t + 0.0006285t^2$$
.

The thermometers used were:—1. One divided into degrees, each of which was 3.5

* Philosophical Magazine for February 1861.

millims. long. With very little practice the temperature could be read off to 0°·1 C. with accuracy. This thermometer was calibrated by ourselves, and afterwards compared with a normal thermometer from Kew Observatory, for which we were indebted to the kindness of Mr. Balfour Stewart. The corrected readings of our thermometer agreed perfectly with those of the Kew thermometer. 2. A normal thermometer from Messrs. Negretti and Zambra, divided into 0°·2 C. This was compared with the Kew thermometer and found to be correct. The boiling- and freezing-points of the thermometers were taken at intervals, and the necessary corrections made.

As the light in the room where the experiments were made came from above, and as the thermometers lay horizontally in the trough, by placing the eye in a position so that the division on the thermometer covered its reflexion on the column of mercury, all error of parallax was avoided. The thermometers were always read off with the help of the magnifying glass A through the oil in the glass tube d, so that the whole of the column of mercury had very nearly the temperature of the bath.

The normal wires were made of annealed german silver, and their resistances determined by comparing them with a hard-drawn wire of the gold-silver alloy*. They were soldered to two thick copper wires, varnished, and when used placed in the cylinder G, filled with oil, in which a thermometer hung. The temperature of the oil was taken immediately after each observation, and the conducting power of the normal wire corrected by the use of the formula

$$\lambda = 7.803 - 0.0034619t + 0.0000003951t^2$$

which was found by the determination of the conducting powers, at different temperatures, of a piece of wire from the same coil as that from which the normal wires were cut. In this paper we have taken as unit the conducting power of a hard-drawn silver wire at 0° C.=100 (that of the hard-drawn gold-silver alloy at 0° being =15.03), in order to be able to compare at sight the present determinations with those made by one of us a short time ago †.

Before beginning a series, as already stated, all the ends of the wires dipping in the mercury-cups were re-amalgamated, and the zero-point of the scale redetermined. The current from the cell D was only allowed to pass through the apparatus for a second or two at a time, for fear of heating the wires, &c.

From 0° to 100° seven intervals were chosen at which observations were made, viz. 12°, 25°, 40°, 55°, 70°, 85°, 100°. With a little practice the flames of the 6-Bunsen burner could be regulated so as to come within a degree or two of the above temperatures. For about five minutes before, and whilst making the observations, the oil in the trough was stirred, one observer being at the trough whilst the other determined the resistances. Four observations at each interval were generally made on heating the wire to 100°, and again four at each interval on cooling (where this was not the case it will be mentioned with the series).

- * Philosophical Magazine, February 1861.
- † Philosophical Transactions, 1858 and 1860.

To save space, the mean only of the eight observations will be given, as otherwise the number of figures would be very great. Table I. may be taken as a fair example of the results obtained. The formulæ from which the conducting powers have been calculated is

$$\lambda = x + yt + zt^2,$$

where λ is the conducting power at t° C., x the conducting power at 0° , and y and z constants. The values for x, y, and z were deduced from the mean of the observations by the method of least squares.

We will now proceed to the experiments made with each metal, making at the same time a few remarks on their purification, &c., and then see what general laws and conclusions we may draw from the results obtained.

Silver:

Purified by precipitating nitrate of silver with hydrochloric acid, and reducing the washed chloride with pure carbonate of sodium. Wires 1, 2, and 3 were of different preparations. Table II. gives the results obtained with these wires.

	First	wire.	Second	wire.	Third wire.			
	Hard drawn.	Annealed.	Hard drawn.	Annealed.	Hard drawn.	Annealed.		
Length Diameter	1546 millims. 0·462 millim.	1535 millims. 0·462 millim.			1962 millims. 0·448 millim.			
Conducting power found before heating the hard-drawn wires	97.645 at 15.4	Reduced to 0°. 103.528	95·112 at 16·0		94·053 at 16·0	Reduced to 0°. 99.800		
after being kept at 100° for 1 day Ditto, for 2 days Ditto, for 3 days Ditto, for 4 days Ditto, for 5 days Ditto, for 6 days	98·138 at 16·2 98·913 at 15·6 99·837 at 16·0 99·212 at 18·4 99·586 at 17·4	104·951 106·091 106·377	96·618 at 15·6 101·544 at 16·8 102·237 at 16·0 101·427 at 19·2 101·750 at 18·6	108·303* 108·714 109·162	95.241 at 15.4 96.337 at 16.0 96.671 at 17.6 97.917 at 15.6 97.669 at 17.4 97.322 at 18.2	102·223 103·178 103·747 104·168		

TABLE II.

^{*} During the day the temperature of the oil increased, by mistake, to 130°.

	First wire, hard drawn.				Second wir	e, hard drav	7n.	Third wire, hard drawn.				
T.	Conducti	ng power.	Difference.	Difference. T.		Conducting power.			Conducti	ng power.	Difference.	
1.	Observed.	Calculated.		1.	Observed.	Calculated.	Dinerence.	T.	Observed.	Calculated.	1	
11.00	102.238	102-272	-0.034	12.20	103-927	103-927	0.000	9 ⋅60	100.534	100.546	-0.012	
26.17	96.710	96.645	+0.065	23.70	99.520	99.523	-0.003	23.90	95.452	95.437	+0.015	
38.25	92.490	92.505	-0.015	41.70	93.224	93.236	-0.012	38.95	90.476	90.507	-0.031	
55.40	87.130	87.149	-0.019	56.20	88.703	88.708	-0.005	56.00	85.513	85.478	+0.035	
68.85	83.389	83.374	+0.015	68.90	85.142	85.137	+0.005	68.15	82.244	82.252	-0.008	
84.00	79.540	79.572	-0.032	85.45	81.078	81.036	+0.042	84.47	78.393	78.391	+0.002	
101.30	75.831	75.813	+0.018	99:20	78.073	78.103	-0.030	98.60	75.477	75.484	-0.007	
	First win	e, annealed.			Second wire, annealed.				Third wire, annealed.			
1 η30	103-391	103.404	-0.013	8 ∙00	106.447	106.426	+0.021	°9.25	102.543	102.461	+0.082	
24.25	98.589	98.576	+0.013	24.35	99.968	99.990	-0.022	25.55	96.371	96.495	-6.124	
41.85	92.520	92.530	-0.010	38.05	95.051	95.077	-0.026	40.10	91.589	91.630	-0.041	
56.45	88.006	87.965	+0.041	55.17	89.554	89.554	0.000	55.17	87.055	87.047	+0.008	
67.75	84.670	84.714	-0.044	68.22	85.847	85.803	+0.044	68.55	83.483	83.367	+0.116	
83.65	80.562	80.554	+0.008	83.62	81.882	81.888	-0.006	83.57	79.667	79.674	-0.007	
98.80	77.046	77.042	+0.004	100.00	78.319	78.331	-0.012	100.00	76.124	76.163	-0.039	

The formulæ deduced from the observations, from which the conducting powers were calculated, were—

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For first wire (hard drawn) . \lambda = 106.651 - 0.40948t + 0.0010370t^2. For first wire (annealed) . \lambda = 107.880 - 0.40698t + 0.0009601t^2. For second wire (hard drawn) \lambda = 108.928 - 0.42389t + 0.0011407t^2. For second wire (annealed) . \lambda = 109.802 - 0.43138t + 0.0011667t^2. For third wire (hard drawn) . \lambda = 104.209 - 0.39124t + 0.0010133t^2. For third wire (annealed) . \lambda = 106.088 - 0.40160t + 0.0010235t^2.
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From the above Table it will be seen that, after heating a silver wire to 100°C. for some days, its conducting power is increased almost to the same extent as if it had been annealed, and that wires 1 and 2 were not completely hard drawn. On comparing the difference in the conducting powers produced by annealing the wires, we find for wire 3 it is only 6 per cent., whereas for wire 2 it is almost 10 per cent., taking the conducting power of the hard-drawn silver wire =100. In a former research* this difference was found to be—

				_	Reduced to 0°.
1.	Hard drawn			95.28 at $1\mathring{4}.0$	100.47
	Annealed .		•	103·98 at 14·8	109.98
2.	Hard drawn			95·36 at 14·6	100.78
	Annealed .			103·33 at 14·6	$109 \cdot 20$

^{*} Philosophical Transactions, 1860.

These values have been reduced by using a formula which is the mean of the six deduced from the experiments; for although there is a difference in the formula obtained for the annealed and hard-drawn (or rather partially annealed) wires, yet it is so small that they may be considered the same, more especially as the difference between the one obtained for the different wires is far greater. Taking the mean of the above values, and assuming the influence of temperature on the conducting power of hard-drawn and annealed wires to be the same, we find the following formulæ:—

For hard-drawn wires $\lambda = 100.00 - 0.38287t + 0.0009848t^2$. For annealed wires $\lambda = 108.574 - 0.41570t + 0.0010624t^2$.

Copper.

Wires 1 and 2 were of the same piece of electrotype copper prepared for us by Dr. H. Müller at Messrs. De la Rue and Co.'s. Wire 3 was cut off a piece of commercial electrotype copper from the same source. Table III. shows the results obtained with these wires.

TABLE III.

	First	wire.	Second	l wire.	Third wire.				
	Hard drawn.	Annealed.	Hard drawn.	Annealed.	Hard drawn.	Annealed.			
Length Diameter			1753 millims. 0·598 millim.						
Conducting power found before heating thehard-drawn wires	95·672 at 10·6	Reduced to 0°. 99·526	94·355 at 15·0	Reduced to 0°. 100.021	92•568 at 2 0 •6	Reduced to 0°. 100·327			
Ditto, for 2 days Ditto, for 3 days Ditto, for 4 days Ditto, for 5 days Ditto, for 6 days Ditto, for 7 days	96·324 at 9·9 96·750 at 11·8 96·914 at 12·2 97·950 at 9·8 98·437 at 8·7	101·097 101·418 101·671 101·682	94.965 at 13.2 94.880 at 14.2 94.501 at 15.9 94.153 at 17.2 95.570 at 14.4 94.327 at 18.2 96.575 at 12.7	100·268 100·524 100·656 101·074	93.263 at 19.0 93.720 at 18.0 93.434 at 19.0 93.278 at 19.6 92.865 at 20.6 92.738 at 21.1	100·645 100·708			

	First wire, hard drawn.				Second wire, hard drawn.				Third wire, hard drawn.			
T.	Conducti	ng power.	Difference.	Difference. T.		ng power.	Difference.	T.	Conducti	ng power.	Difference.	
1.	Observed.	Calculated.		1.	Observed.	Calculated.	1	т.	Observed.	Calculated.	1	
16.86 29.88 51.03 69.52 83.77 98.60	95·473 91·063 84·235 78·997 75·413 71·829	95·467 91·002 84·315 79·044 75·347 71·838	+0.006 +0.061 -0.080 -0.047 +0.066 -0.009	19·17 30·95 48·53 69·22 83·77 99·00	94·359 90·187 84·518 78·640 75·015 71·532	94·334 90·208 84·544 78·634 74·968 71·562	+0.025 -0.021 -0.026 +0.006 +0.047 -0.030	12.65 25.61 39.52 53.92 69.90 84.87 99.92	95·769 91·061 86·415 82·069 77·798 74·172 70·951	95·739 91·076 86·456 82·090 77·741 74·142 70·987	+0.030 -0.015 -0.041 -0.021 +0.057 +0.030 -0.036	
	First wi	re, annealed.		Second wire, annealed.				Third wire, annealed.				
1 ⁵ ·00 29·63 50·22 69·60 83·42 99·39	95·535 91·291 84·687 79·223 75·636 71·891	95·567 91·239 84·726 79·209 75·638 71·893	-0.032 +0.052 -0.039 +0.014 -0.002 -0.002	18.96 31.86 52.05 70.27 83.81 99.57	94·987 90·424 83·974 78·836 75·428 71·757	94·959 90·449 84·003 78·829 75·377 71·784	+0.028 -0.025 -0.029 +0.007 +0.051 -0.027	13·45 26·15 39·35 55·50 69·90 84·67 99·05	96·954 92·246 87·727 82·675 78·742 75·047 71·766	96·934 92·260 87·753 82·722 78·686 74·988 71·816	$\begin{array}{c} +0.020 \\ -0.014 \\ -0.026 \\ -0.047 \\ +0.056 \\ +0.059 \\ -0.050 \end{array}$	

The formulæ deduced from the observations, from which the conducting powers were calculated, were—

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For first wire (hard drawn) . . . \lambda = 101 \cdot 645 - 0 \cdot 37963t + 0 \cdot 0007844t^2. For first wire (annealed) . . . \lambda = 101 \cdot 791 - 0 \cdot 37959t + 0 \cdot 0007921t^2. For second wire (hard drawn) . . \lambda = 101 \cdot 614 - 0 \cdot 39806t + 0 \cdot 0009546t^2. For third wire (annealed) . . . \lambda = 102 \cdot 143 - 0 \cdot 39629t + 0 \cdot 0009179t^2. For third wire (hard drawn) . . . \lambda = 100 \cdot 620 - 0 \cdot 39885t + 0 \cdot 0010236t^2. For third wire (annealed) . . . \lambda = 102 \cdot 243 - 0 \cdot 40850t + 0 \cdot 0010228t^2.
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The observations made with wires 1 and 2 were as follows: two at each interval on heating and two on cooling; again, two on heating and two on cooling, as shown in Table I.

On looking at the above, we observe that wire 1, after having been kept at 100° for several days, increased in conducting power almost to the same extent as if it had been annealed, wire 2 partially so, and wire 3 hardly at all. The annealing took place in a glass tube heated with a 4-Bunsen burner, whilst a current of hydrogen passed through it. Here, again, as in the case of the silver wire, we may assume that the formulæ of the hard-drawn and annealed copper wires are the same. In a former research* pure copper was found to conduct—

	0	Reduced to 0°.
1.	93.00 at 18.6	99.877
2.	93.46 at 20.2	100.980
3.	92.02 at 18.4	99.824
4.	92.76 at 19.3	99.886
5.	92.99 at 17.5	99.453

^{*} Philosophical Transactions, 1860.

The difference found between the conducting powers of hard-drawn and annealed wires was—

c	TT 1 1						95·31 at $1\mathring{1}$ ·0	Reduced to 0°.
0.	Hard drawn	•	•	•	•	•	99.91 at 11.0	$99 \cdot 435$
	Annealed.	•	•		•		97·83 at 11·0	$102 \cdot 065$
7.	Hard drawn						95·72 at 11·0	99.864
	Annealed.		•		٠.		98.02 at 11.0	$102 \cdot 263$

These values have been reduced to 0° as follows: take for instance the first, 93.00 at 18°.6. The mean of the six formulæ obtained for copper (see Table XV.) is

$$\lambda = 100 - 0.38701t + 0.0009009t^2$$
;

and calculating the conducting power for 18°.6 by this formula, we find it equal to 93.114.

Now
$$\frac{93.00}{93.114} = 0.99877;$$

and if all the terms of the above formula be multiplied by this number, we deduce a formula by which the above value can be reduced. All the reductions given in this paper of former determinations were made in this manner, using the formulæ given in Table XV. The reductions to 0° in the Tables were made in a like manner, the only difference being that the formulæ found for the respective wires were used instead of the mean. Taking the mean of all the values found for copper, and using the mean for the formulæ given in Table XV., we find as the formula for correction of the conducting power for temperature of

A hard-drawn wire
$$\lambda = 99.947 - 0.38681t + 0.0009004t^2$$

An annealed wire $\lambda = 102.213 - 0.39557t + 0.0009208t^2$.

The values given as first term in the formulæ were found as follows: on referring to the paper* from which the conducting powers of copper were taken, it will be seen that each of them is the mean of three determinations. The reduced values therefore of 1 to 5, the mean of 6, hard drawn, and 7, hard drawn, and the mean of the first determinations of the three wires given in Table III., were added together, and the mean taken as the conducting power of a hard-drawn copper wire at 0° C. For the annealed, the per-centage differences of the values of 6, hard drawn and annealed, 7, ditto, and of the first determinations of the three wires in Table III. and the annealed ones, were added together, and the mean added to the value found for the hard-drawn wire (as a percentage amount). All the formulæ given as end-result with each metal have been constructed in this manner.

Gold.

Purified as described in the Philosophical Transactions, 1860, p. 175. Wires 1, 2, and 3 were of different preparations. The results obtained with these wires are given in Table IV.

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A TOT TO	1:3/
TABLE	. v .

	First	wire.	Second	wire.	Third wire.			
	Hard drawn.	Annealed.	Hard drawn.	Annealed.	Hard drawn.	Annealed.		
Length Diameter	2214 millims. 0·759 millim.	2200 millims. 0·759 millim.	837 millims. 0·467 millim.		759·5 millims. 0·434 millim.	•		
Conducting power found before heating the hard-drawn wires	73·239 at 13 · 2	Reduced to 0°. 76.821	72·550 at 15·1	Reduced to 0°. 76.561	67·530 at 36·8	Reduced to 0°. 77•229		
100° for 1 day Ditto, for 2 days Ditto, for 3 days Ditto, for 4 days	72.746 at 15.2 72.751 at 15.1		73·359 at 12·6	76·733 	71.868 at 19.4 71.854 at 20.1 72.191 at 19.0 72.396 at 18.0	77 · 405 77·457		

The means of the conducting powers found for each of the following temperatures were—

	First wire,	hard drawn	1.	Second wire, hard drawn.				Third wire, hard drawn.			
	Conducti	ng power.	TD: 00	_ m	Conducti	ng power.			Conducti	ng power.	D:00
Т.	Observed.	Calculated.	Difference.	т.	Observed.	Calculated.	Difference.	T.	Observed.	Calculated.	Difference.
15.95 30.76 50.06 69.75 84.31 99.27	72·567 68·798 64·392 60·397 57·742 55·248	72·536 68·828 64·410 60·385 57·722 55·263	$\begin{array}{c} + 0.031 \\ - 0.030 \\ - 0.018 \\ + 0.012 \\ + 0.020 \\ - 0.015 \end{array}$	13·36 24·79 40·80 55·65 69·52 84·12 100·00	73·222 70·329 66·515 63·306 60·528 57·905 55·203	73·212 70·325 66·544 63·312 60·531 57·854 55·232	+0.010 +0.004 -0.029 -0.006 -0.003 +0.051 -0.029	12·44 23·27 39·42 55·47 70·56 84·79 99·00	73.854 70.965 67.002 63.441 60.455 57.904 55.635	73·841 70·975 67·013 63·448 60·435 57·893 55·647	$\begin{array}{c} +0.013\\ -0.010\\ -0.011\\ -0.007\\ +0.020\\ +0.011\\ -0.012 \end{array}$
	First wir	e, annealed.		Third wire, annealed.							
14.92 30.05 48.87 69.90 82.82 99.62	70.039 65.575 61.220 58.811	73·992 70·068 65·611 61·191 58·768 55·948	+0.028 -0.029 -0.036 +0.029 +0.043 -0.033	minatio		e was torn	the deter- away from		74·327 71·067 67·582 64·645 61·229 58·422 56·029	74·293 71·095 67·621 64·628 61·220 58·388 56·056	+0.034 -0.028 -0.039 +0.017 +0.009 +0.034 -0.027

The formulæ deduced from the observations, from which the conducting powers were calculated, were—

For first wire (hard drawn). . $\lambda = 76.838 - 0.27973t + 0.0006285t^2$.

For first wire (annealed) . . . $\lambda = 78.161 - 0.28935t + 0.0006664t^2$.

For second wire (hard drawn) . $\lambda = 76.786 - 0.27549t + 0.0005995t^2$.

For third wire (hard drawn . . $\lambda = 77.343 - 0.29043t + 0.0007200t^2$.

For third wire (annealed) . . . $\lambda = 78.231 - 0.28849t + 0.0006564t^2$.

The observations made with wire 1 (hard drawn) are given in Table I., those of the same wire (annealed) were made in the same manner.

Here we find no permanent change in conducting power with wire 1, after being kept at 100° for several days, and only a very slight increase with wires 2 and 3. The formulæ for the hard-drawn and annealed wires agree so closely that they may also, as with silver and copper, be considered the same.

In the paper just alluded to, the conducting power of pure gold was found—

1.	72·68 at 19°·3	Reduced to 0° . 77.966
2.	73.08 at 23.3	79.524
3.	73·27 at 13·8	77.053
4.	73·99 at 15·1	78.178

The difference between hard-drawn and annealed wires was—

5.	Hard drawn	•	•	•	•	•	74·20 at 14·8	Reduced to 0°. 78.313
	Annealed .						75.53 at 15.2	79.833
6.	Hard drawn	•	•		•		73·78 at 15·5	78.067
	Annealed .		٠.				75·18 at 15·8	79.635

Taking the mean of the values as with copper, the following formulæ were deduced for the correction of conducting power for temperature:—

For hard-drawn wires $\lambda = 77.964 - 0.28648t + 0.0006582t^2$. For annealed wires $\lambda = 79.327 - 0.29149t + 0.0006697t^2$.

Zinc.

Zinc free of arsenic was purified by distillation. All pressed wires. In Table V. the results obtained are given.

TABLE V.

	First wire.	Second wire.	Third wire.		
Length Diameter	502•2 millims. 0•588 millim.	394 millims. 0·513 millim.	372 millims. 0·519 millim.		
Conducting power found before heating the wires Ditto, after being kept at 100° for	Reduced to 0 26.744 at 23.1 29.093	e. Reduced to 0°. 26.903 at 18.5 28.836	Reduced to 0°. 26.835 at 18.0 28.639		
1 dayDitto, for 2 days	26·695 at 23·7 29·103	27.081 at 17.5 28.919 26.980 at 18.5 28.919	26.784 at 18.5 28.636 26.885 at 17.4 28.632		

	Conducti	ng power.	T: m		Conducti	ng power.	- I		Conducti	ng power.	7.0
T.	Observed.	Calculated.	Difference.	T.	Observed.	Calculated.	Difference.	T.	Observed.	Calculated.	Difference.
11.60 24.24 41.33 55.08 70.27 82.01 98.07	27·915 26·639 25·077 23·925 22·757 21·924 20·865	27·902 26·653 25·086 23·926 22·747 21·912 20·875	+0.013 -0.014 -0.009 -0.001 +0.010 +0.012 -0.010	11°20 26·12 39·55 54·18 72·32 85·77 100·23	27.706 26.187 24.951 23.719 22.330 21.407 20.540	27.687 26.199 24.959 23.716 22.330 21.414 20.534	+0.019 -0.012 -0.008 +0.003 -0.000 -0.007 +0.006	11°·16 25·96 40·10 56·85 71·73 85·40 98·95	27.518 26.088 24.812 23.423 22.306 21.348 20.462	27·513 26·090 24·820 23·428 22·295 21·339 20·472	+0.005 -0.002 -0.008 -0.005 +0.011 +0.009 -0.010

The formulæ deduced from the observations, from which the conducting powers were calculated, were—

For first wire . . $\lambda = 29 \cdot 114 - 0 \cdot 10727t + 0 \cdot 0002372t^2$. For second wire . $\lambda = 28 \cdot 881 - 0 \cdot 10949t + 0 \cdot 0002616t^2$. For third wire . . $\lambda = 28 \cdot 649 - 0 \cdot 10424t + 0 \cdot 0002182t^2$.

No permanent alteration in the conducting power takes place after heating the wires for several days to 100°.

The value formerly found for the conducting power of zinc (precipitated galvanoplastically, fused and pressed) was—

Reduced to 0° . 29.220.

Treating these values as before, we find the formula for zinc to be

$$\lambda = 29.022 - 0.10752t + 0.0002401t^2$$
.

Cadmium.

The metal was purified as described in the Philosophical Transactions, 1860, p. 177. The wires were pressed. Table VI. shows the results.

TABLE VI.

	First wire.	Second wire.	Third wire.
Length Diameter	625 millims.	559 millims.	439 millims.
	0·641 millim.	0.678 millim.	0·684 millim.

	Conducti	ng power.]		Conducti	ng power.			Conducti	ng power.	D: 00
т.	Observed.	Calculated.	Difference.	T.	Observed.	Calculated.	Difference.	T.	Observed.	Calculated.	Difference.
8·87 20·75 34·47 49·38 63·39 77·74 93·55	21·241 20·138 19·188 18·292	23·329 22·338 21·255 20·150 19·186 18·268 17·339	$\begin{array}{c} -0.002 \\ +0.013 \\ -0.014 \\ -0.012 \\ +0.002 \\ +0.024 \\ -0.014 \end{array}$	8·89 21·59 36·37 48·52 62·90 80·00	23·374 22·280 21·075 20·157 19·171 18·109	23·400 22·270 21·059 20·146 19·162 18·131	-0.026 +0.010 +0.016 +0.011 +0.009 -0.022	14.60 22.05 39.65 54.45 68.10 81.20 89.90	21.849 21.318 20.072 19.065 18.179 17.393 16.896	21.859 21.310 20.061 19.067 18.194 17.397 16.888	$\begin{array}{c} -0.010 \\ +0.008 \\ +0.011 \\ -0.002 \\ -0.015 \\ -0.004 \\ +0.008 \end{array}$

The formulæ deduced from the observations, from which the conducting powers were calculated, were—

For first wire . $\lambda = 24 \cdot 100 - 0 \cdot 088554t + 0 \cdot 0001740t^2$. For second wire . $\lambda = 24 \cdot 240 - 0 \cdot 096753t + 0 \cdot 0002548t^2$. For third wire . $\lambda = 24 \cdot 974 - 0 \cdot 078004t + 0 \cdot 0001147t^2$.

The values obtained for the alteration in the conducting power of these wires after heating them for several days to 100°, have unfortunately been lost. It may, however, be stated that the differences were very small, and that there was a loss in conducting power.

The conducting power of cadmium was found in the paper already referred to—

22·10 at 18·8 Reduced to 0°. 23·678.

Deducing the formula for cadmium in the manner before described, we find

 $\lambda = 23.725 - 0.087476t + 0.0001797t^2$.

Pure cadmium, when heated to about 80°, becomes exceedingly brittle, in fact it may be powdered in a hot mortar with great ease. We should not have been able to carry out the determinations if the wires had not been varnished, as the movement of the oil by the stirrer would have caused them to fall to pieces. It is worthy of remark that this change in the molecular arrangement of the wires does not make itself apparent in the conducting power to any very marked extent.

Tin.

Purified by dissolving commercial tin in nitric acid, and reducing the washed oxide by heating it with lampblack. Pressed wires were used. Table VII. gives the results.

	First win	re.		Second	wire.	rain e	Third v	vire.
LengthDiameter	279 mil 0•559 mil		0	375 m ·634 m		315 millims. 0·729 millim.		
Conducting power found before heating the wires Ditto, after being kept at 100° for 1		Reduced to 0°·	11.532	at 18•1	Reduced to 0°.		at 18.2	Reduced to 0° 13·108
day Ditto, for 2 days	10.852 at 27.0	11·926 11·956 11·980	11·442 11·448 11·444	at 18.6	6 12.257	12·291 12·296	at 18·4 at 18·4	

TABLE VII.

T.	Conducti	ng power.	Difference	T.	Conducti	ng power.	Difference.	T.	Conducti	ng power.	Difference.
4.	Observed.	Observed. Calculated.		1.	Observed.	Calculated.			Observed.	Calculated.	
25.27	11·4110 10·9320 10·3570 9·8498 9·2980 8·9033 8·3937	11·4202 10·9246 10·3436 9·8558 9·3046 8·9078 8·3881	$\begin{array}{c} -0.0092 \\ +0.0074 \\ +0.0134 \\ -0.0060 \\ -0.0066 \\ -0.0045 \\ +0.0056 \end{array}$		11.7144 11.1287 10.5805 10.0451 9.4883 9.0102 8.6158	11·7227 11·1153 10·5732 10·0526 9·4961 9·0127 8·6096	-0.0083 +0.0134 +0.0073 -0.0075 -0.0078 -0.0025 +0.0062	10.00 26.54 39.52 56.27 70.30 85.72 97.30	12.649 11.944 11.408 10.717 10.189 9.654 9.279	12.660 11.934 11.391 10.727 10.202 9.657 9.270	-0.011 +0.010 +0.017 -0.010 -0.013 -0.003 +0.009

The formulæ deduced from the observations, and from which the conducting powers were calculated, were—

For first wire $\lambda = 11.9613 - 0.042902t + 0.00007422t^2$. For second wire $\lambda = 12.2419 - 0.044965t + 0.00008213t^2$.

For third wire $\lambda = 13.1186 - 0.046561t + 0.00007206t^2$.

We see from the results that wires 1 and 2 decrease to a small extent in conducting power, whereas wire 3 increases slightly after being heated to 100°.

The conducting power of tin was found—

11.45 at $2\mathring{1}.0$

Reduced to 0° . 12.351:

and calculating the formula of tin as before, we find

 $\lambda = 12.366 - 0.044554t + 0.00007588t^2$

Lead.

Purified by reducing by heat the twice recrystallized acetate. Wires 1 and 2 were pressed; wire 3 drawn. No permanent alteration in the conducting power of the wires was observed after they had been kept at 100° for two days. Table VIII. shows the results.

TABLE VIII.

	First wire.	Second wire.	Third wire.
Length	0.000 1111	453 millims.	389 millims.
Diameter		0.698 millim.	0·959 millim.

	Conducti	ng power.	Difference.	T.	Conducti	ng power.	Difference.	T.	Conducti	ng power.	Difference.
T.	Observed:	Calculated.		1.	Observed.	Calculated.			Observed.	Calculated.	
14.5. 25.4. 40.3 54.8 70.3 84.5 99.3	7.6129 7.2036 6.8423 6.4881 6.1964	7·9336 7·6152 7·2071 6·8420 6·4863 6·1929 5·9189	+0.0029 -0.0023 -0.0035 +0.0003 +0.0018 +0.0035 -0.0030	lt	7·8685 7·5336 7·1405 6·7789 6·4392 6·1189 5·8388	7.8653 7.5392 7.1397 6.7775 6.4370 6.1218 5.8381	+0.0032 -0.0056 +0.0008 +0.0014 +0.0022 -0.0029 +0.0007	12.40 26.20 39.60 54.60 69.70 84.40 98.85	7.9038 7.4967 7.1309 6.7565 6.4205 6.1250 5.8642	7.9022 7.4968 7.1324 6.7585 6.4187 6.1229 5.8658	+0.0016 -0.0001 -0.0015 -0.0020 +0.0018 +0.0021 -0.0016

The formulæ deduced from the observations, from which the conducting powers were calculated, were—

For first wire . $\lambda = 8.3882 - 0.032346t + 0.00007540t^2$. For second wire $\lambda = 8.3147 - 0.032055t + 0.00007307t^2$. For third wire . $\lambda = 8.2925 - 0.032468t + 0.00008011t^2$.

The value found for the conducting power of lead was

7.77 at 17.3 Reduced to 0°. 8.304.

Treating the mean of the values as above, the formula is

 $\lambda = 8.318 - 0.032237t + 0.00007608t^2$.

Arsenic.

Purified by sublimation. Small bars were cut from a comparatively solid piece and soldered to two copper wires; on account of the extreme brittleness of arsenic, the bars were placed in glass tubes closed at the ends with gypsum, through which the copper wires passed. As these were dried in a water-bath for several days, no permanent alteration of the conducting power of the bars was found after being heated in the oil-bath for two days. The values found for the conducting power of arsenic agree as well as could be expected, considering the bars were made by hand, and the metal somewhat porous. The difficulty of obtaining bars of metal of sufficient length is so great that we have been contented with two series. These are given in Table IX.

TABLE IX.

	First bar.	Second bar.
Length Diameter	0.00	55·5 millims. 1·01 millim.

The means of the conducting powers found for each of the following temperatures were—

T.	Conducti	ng power.	Difference.	T.	Conducti	ing power.	Difference.	
1.	Observed.	Calculated.	Difference.	1.	Observed.	Calculated.	Dinerence.	
14·20 25·30 37·80 55·00 70·00 85·30 101·00	5·0203 4·8007 4·5710 4·2854 4·0767 3·8810 3·7005	5.0180 4.8008 4.5736 4.2906 4.0722 3.8764 3.7041	+0.0023 -0.0001 -0.0026 -0.0052 +0.0045 +0.0046 -0.0036	13.50 24.50 40.15 55.55 69.90 82.50 99.80	4·0051 3·8371 3·6367 3·4447 3·2559 3·1144 2·9485	4·0037 3·8450 3·6311 3·4341 3·2628 3·1221 2·9435	+0.0014 -0.0079 +0.0056 +0.0106 -0.0069 -0.0077 +0.0050	

The formulæ deduced from the observations, and from which the conducting powers were calculated, were—

For first bar $\lambda = 5.3168 - 0.021874t + 0.00005848t^2$. For second bar $\lambda = 4.2078 - 0.015506t + 0.00002843t^2$. Taking the mean of the conducting powers at 0°, we deduce the formula for the correction of conducting power for temperature to be

 $\lambda = 4.7623 - 0.018571t + 0.00004228t^2$.

Antimony.

Purified by twice recrystallizing commercially pure tartrate of antimony and potassium, reducing by heat and re-fusing with antimonic acid. As antimony is so very brittle, it was not possible to manipulate with it in form of wire, it was therefore fused in the bowl of a tobacco-pipe, and when liquid allowed to run into the stem. After breaking off the bowl, the ends of the pipe were made so hot that the metal melted, and clean copper wires were pushed into the liquid metal, which on solidifying held them fast. The free ends of the copper wires were then soldered to the thick ones in the Unfortunately in each case the copper wires in the pipe-stem became loose trough. after heating for two or three days, and had to be therefore resoldered, so that no reliable determinations could be made as to the effect of heating to 100° for several days on the conducting power. It may be stated that the three wires lost in conducting power; but to what extent, we are of course not in a position to say. As the diameter of the pipe-stem could not be accurately determined, and as it could not be ascertained whether there were cavities in the wires (caused by contraction on cooling and crystallization) or not, the first observed conducting power was taken equal to 100. Table X. shows the results.

Table X.

The means of the conducting powers found for each of the following temperatures were—

	Firs	st wire.			Seco	nd wire.		Third wire.			
T.	Conducting power. Difference.				Conducting power.		т.	Conducting power.		Difference	
1.	Observed.	Calculated.		1.	Observed.	Calculated.	Difference.	1.	Observed.	Calculated.	Difference.
10.00 26.35 40.40 54.55 70.65 83.50 99.40	100·000 94·062 88·982 84·633 80·126 77·071 73·430	100·052 93·910 89·089 84·664 80·152 76·953 73·484	-0.052 +0.152 -0.107 -0.031 -0.026 +0.118 -0.054	8·40 25·60 42·45 57·80 69·45 86·85 101·25	100·000 93·947 88·139 83·707 80·691 76·138 72·922	99·999 93·850 88·329 83·731 80·517 76·159 72·953	+0.001 +0.097 -0.190 -0.024 +0.174 -0.021 -0.031	13.80 22.30 38.65 53.50 69.65 84.45 98.80	100·000 96·378 90·552 85·671 81·118 77·480 74·448	99·901 96·514 90·527 85·692 81·082 77·454 74·480	+0.099 -0.136 $+0.025$ -0.021 $+0.036$ $+0.026$ -0.032

The formulæ deduced from the observations, from which the conducting powers were calculated, were—

For first wire $\lambda = 104.095 - 0.41487t + 0.0010755t^2$.

For second wire $\lambda = 103 \cdot 190 - 0.38721t + 0.0008748t^2$.

For third wire $\lambda = 105.801 - 0.44541t + 0.0012995t^2$.

MDCCCLXII.

The observed conducting powers in this and the foregoing Table do not agree so well with the calculated as the others, on account of the temperature of the bath never being exactly the same as that of the wire; for in the one case the heat had to traverse the glass tube filled with air, in the other the thickness of the pipe-stem, before reaching the metal.

The conducting power of antimony was found equal to

Reduced to 0° . 4.6172

Using this value as before described, we obtain a formula for antimony where $\lambda = 4.6172 - 0.018389t + 0.00004788t^2$.

Bismuth.

Purified by reducing the basic nitrate of bismuth with lampblack. Table XI. gives the results. The wires were pressed.

	First wire.			Second	wire.	Third wire.		
Length Diameter	117 mi 0·596 mi	1	121•4 m 0•596 m		42.5 millims. 0.217 millim.			
Conducting power found before heating the wires Conducting power after being kept at	1·1787 at 16·6	Reduced to 0°. 1.2517	1.1036	at 18.8	Reduced to 0°.	1.2215	at 16.6	Reduced to 0°. 1.2951
100° for 1 day Ditto, for 2 days Ditto, for 3 days	1.3595 at 18.2	1·4494 1·4521 1·4541	1.3121	at 19.0 at 19.0 at 19.9	1.4006	1.3709	at 17.8 at 17.6 at 17.9	1.4587

TABLE XI.

The means of the conducting powers found for each of the following temperatures were—

T.	Conducti	ng power.	Difference.	т.	Conducti	ng power.	Difference.	T.	Conducti	ng power.	Difference.	
1.	Observed.	Calculated.		1.	Observed.	Calculated.			Observed.	Calculated.	Z IZZ CHOO.	
9.20 26.15 39.50 57.25 68.95 84.35 96.35	1·4059 1·3226 1·2609 1·1863 1·1397 1·0833 1·0428	1·4058 1·3226 1·2614 1·1858 1·1397 1·0833 1·0429	+0.0001 0.0000 -0.0005 +0.0005 0.0000 0.0000	8.60 24.00 38.75 55.30 68.90 84.00 95.90	1:3654 1:2909 1:2297 1:1591 1:1058 1:0478 1:0036	1·3641 1·2935 1·2287 1·1593 1·1050 1·0474 1·0042	+0.0013 -0.0026 +0.0010 -0.0002 +0.0008 +0.0004 -0.0006	9.40 25.65 43.05 57.45 71.60 88.60	1·4129 1·3329 1·2551 1·1913 1·1315 1·0671	1.4128 1.3339 1.2538 1.1912 1.1328 1.0666	$\begin{array}{c} +0.0001 \\ -0.0010 \\ +0.0013 \\ +0.0001 \\ -0.0013 \\ +0.0005 \end{array}$	

The formulæ deduced from the observations, by which the conducting powers were calculated, were—

For first wire . $\lambda = 1.4535 - 0.0052883t + 0.00001060t^2$. For second wire $\lambda = 1.4049 - 0.0047972t + 0.000006453t^2$. For third wire . $\lambda = 1.4603 - 0.0051286t + 0.000007737t^2$.

From the above we see how bismuth increases in conducting power after being kept at 100° for one day. This increment is so rapid that it may be followed for the first two hours from five to five minutes. Wire 1 altered by one day's heating 16 per cent.; wire 2, 19 per cent.; and wire 3, 12 per cent. Wires 1 and 2 were cut from the same piece.

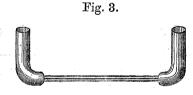
This behaviour explains why the conducting power of bismuth wires varies so much: for in the paper so often here alluded to, the maximum difference between twelve wires was found to be 22 per cent. In pressing the wires the heat applied to the press is never constant; so that, if pressed very warm, wires of high conducting power would probably be the result. The conducting power of bismuth was found equal to

Taking the mean of the values as before, we find the formula for bismuth to be $\lambda = 1.2454 - 0.0043858t + 0.000007134t^2$.

Mercury.

Purified by allowing a solution of subnitrate of mercury to stand over the metal for several weeks, during which time it was often well shaken up with it. The determinations were made in a calibrated thermometer-tube, to the ends of which wide glass tubes (13 to 14 millims. wide) were fused and bent, as shown in fig. 3. Mercury prepared at

different times was used for the determinations. For the experiments, the tube was filled with hot mercury, and its resistance was determined when cold. This was twice repeated; and the resistance being found the same each time, it was assumed that the tube filled in this manner did not contain air-bubbles; this is also proved by the close agree-



ment of the formulæ found in the two cases for the variation of the conducting power at higher temperatures; for if in either case air-bubbles had been present, the formulæ must have differed to a much greater extent, as it can scarcely be assumed that in the two cases the bubbles were equal in bulk. The mercury was connected with the apparatus by amalgamated copper wires (4 to 5 millims. thick). Table XII. shows the results obtained.

TABLE XII.

Length ... = 269 millims. Diameter = 1.424 millim.

The means of the conducting powers found for each of the following temperatures were—

T,	Conducti	ng power.	Difference.	T.	Conducti	Difference.		
1.	Observed.	Observed. Calculated.			Observed.	Calculated.	Dincience.	
0 20·55 40·45 59·82 79·78 99·90	1.6521 1.6276 1.6003 1.5750 1.5465 1.5162	1.6530 1.6272 1.6011 1.5746 1.5462 1.5164	-0.0009 +0.0004 -0.0008 +0.0004 +0.0003 -0.0002	0 20·95 39·92 60·40 80·70 99·30	1.6529 1.6272 1.6010 1.5741 1.5454 1.5174	1.6533 1.6268 1.6018 1.5738 1.5450 1.5177	-0.0004 +0.0004 -0.0008 +0.0003 +0.0004 -0.0003	

The formulæ deduced from the observations, by which the conducting powers were calculated, were—

For the first series . $\lambda = 1.6530 - 0.0012240t - 0.000001434t^2$. For the second series $\lambda = 1.6533 - 0.0012370t - 0.000001297t^3$.

The value found for the conducting power of mercury was

$$^{\circ}$$
 Reduced to 0°. 1.63 at 22.8 1.6588

Taking the mean of the values as before, we find the formula for mercury to be $\lambda = 1.656 - 0.0012326t - 0.000001368t^2$.

Tellurium.

Purified by dissolving the commercial metal in aqua regia, evaporating to dryness with excess of carbonate of sodium, fusing the residue, which was dissolved in water, and nitrate of barium added to precipitate any selenium present. The filtrate was evaporated to dryness with hydrochloric acid in excess, the residue dissolved in water, and precipitated by sulphurous acid.

On account of the low conducting power of tellurium, small bars of about 15 millims. in length and 3-5 millims. in diameter were used for the experiments. Bars I. and II. are of the same preparation. As the bars could not be accurately measured, we have called the first observed conducting power 100 in each case. Table XIII. gives the results.

TABLE XIII.

5	Bar I.			Bar II.	3.1		Bar III	
Conducting power found before heating the bars to 100°.	100 at 16·4	•••••	100	at 15°9	****	100	at 15.6	
Ditto, after being kept at 100° for 1 day	79·145 at 15·4		86 ·50	at 13.0		83·16	at 12.6	
Ditto, for 2 days	45.449 at 16.0	•••••	76·51 70·43	at 13.6 at 16.4	••••	69 ·23 61 ·25	at 14·1 at 16·9	
Ditto, for 3 days Ditto, for 4 days	22.378 at 16.0 16.129 at 15.0		65.68	at 16.6	•••••	54.92	at 17.2	
Ditto, for 5 days	8.068 at 15.2	••••	61.68 56.85	at 16.8 at 17.2		50·69 46·11	at 17.8 at 16.6	
Ditto, for 6 days Ditto, for 7 days	6.989 at 15.0 5.781 at 14.2	*****	54.88	at 16.6	•••••	42.35	at 16.4	
Ditto, for 8 days	4.830 at 15.5 4.621 at 16.8	•••••	51·33 46·27	at 16·1 at 15·6.		$38.64 \\ 35.31$	at 15.8 at 16.2	
Ditto, for 9 days Ditto, for 10 days	4.302 at 15.3	•••••	45.26	at 16.2		33.50	at 16.4	
Ditto, for 11 days	4·181 at 15·0 4·1371 at 16·1	Reduced to 0° . 3.7662	42·10 41·31	at 16.6 at 17.4	•••••	30·97 29·98	at 16.8 at 18.2	
Ditto, for 12 days Ditto, for 13 days	4.0844 at 14.6	- •	39.28	at 16.0 at 17.1	•••••	28·21 26·73	at 15.6 at 16.8	
Ditto, for 14 days Ditto, for 15 days	•••••	•••••	37·72 35·35	at 15.4	•••••	23.68	at 15.4	
Ditto, for 16 days	•••••	•••••	32.23	at 15.6 at 17.0	•••••	19·43 16·65	at 16.0 at 17.6	
Ditto, for 17 days Ditto, for 18 days	•••••		28.11	at 17.6	*****	14.43	at 17.0	
Ditto, for 19 days Ditto, for 20 days			26·25 25·54	at 16.2 at 13.0	*****	12·59 11·68	at 16.4 at 14.4	
Ditto, for 21 days	******		24.12	at 13.4 at 12.8	•••••	10.34	at 13.6 at 13.6	
Ditto, for 22 days Ditto, for 23 days			22.00	at 13.6	•••••	8.64	at 14:1	
Ditto, for 24 days	*******		21.45	at 14·1 at 14·6		7.92	at 13.8 at 14.6	
Ditto, for 25 days Ditto, for 26 days	*******		20.17	at 15.8		6.97	at 14.2	
Ditto, for 27 days	***************************************		19.74	at 16.0 at 13.0		6.66	at 14.8 at 15.8	
Ditto, for 28 days Ditto, for 29 days	******		19.65	at 12·2 3 at 12·0	Reduced to 0° 20·145	6.35	at 15.8 at 12.6	
Ditto, for 30 days Ditto, for 31 days	•••••	•••••		3 at 12.0 3 at 11.9	20.145	6.04	at 12.0	Reduced to 0°
Ditto, for 32 days							0 at 11.8 2 at 12.2	
Ditto, for 33 days	•••••	•••••	1	· · · · · · · · · · · · · · · · · · ·	<u> </u>	1 300	- WU = - 1	1

	Conducti	ng power.			Conducti	ng power.	T): m	m	Conducti	ng power.	Difference.
T.	Observed.	Calculated.	Difference.	T.	Observed.	Calculated.	Difference.	T.	Observed.	Calculated.	
10·40 25·25 38·85 55·10 70·45 83·10 99·40	3·9566 4·5212 5·3940 6·9089 8·8706 11·0316 14·3690	3·9575 4·5240 5·3846 6·9060 8·9019 11·0018 14·3764	-0.0009 -0.0028 +0.0094 +0.0029 -0.0313 +0.0298 -0.0074	11.80 22.40 29.40 29.40 34.60	19.976 19.650 19.477 19.468 19.546 19.689 20.513 21.825 25.096 29.765	19.972 19.660 19.466 19.473 19.496 19.544 19.653 20.464 21.942 25.019 29.784	+0.004 -0.010 +0.011 -0.005 -0.028 +0.002 +0.036 +0.049 -0.117 +0.077 -0.019	4.20 22.80 39.40 53.90 69.50 83.60 98.80	5.6646 6.7456 8.6703 10.8480 14.2472 18.4043 23.3209	5.6797 6.7392 8.5781 10.9426 14.3441 18.2406 23.3769	-0.0151 +0.0064 +0.0922 -0.0946 -0.0969 +0.1637 -0.0560

The formulæ deduced from the observations, and from which the conducting powers were calculated, were—

```
For first bar . . . . . \lambda = 3.7619 + 0.011614t + 0.0006598t^2 + 0.000002994t^3.
```

For second bar to 29.4 . $\lambda = 20.162 - 0.055338t + 0.001085t^2$.

For second bar from 29.4 to $100 \lambda = 20.014 - 0.029569t + 0.00009390t^2 + 0.000010635t^3$.

For third bar $\lambda = 5.5752 + 0.019274t + 0.0013235t^2 + 0.000003088t^3$.

From the above Table we learn that tellurium behaves in a very different manner from the other metals; for it will be seen how very much the conducting power decreases after it has been heated to 100° for some days, and how different is the time required before the conducting power of the different bars becomes constant, or, in other words, until the heating of the bars to 100° causes no further permanent alteration in the conducting power. Bar I. required 13 days; bar II. 32; bar III. 33. The first observed conducting power being taken equal to 100, bar I. is reduced to 4, bar II. to 19·6, and bar III. to 6. If we now look at the determinations of the conducting power at different temperatures of the three bars, we are struck at the great want of concordance in the results. With the first series we observe that the conducting power increases rapidly as the temperature rises; with the second it decreases with the rise of temperature to 29°·4, from which point it increases rapidly, as with bar I.; the third behaves as the first.

Bar I. showed no apparent difference in crystalline structure after being heated; it was thought very probable that the crystalline structure might have been altered by heating, and thus caused the enormous change in conducting power. The three bars, when first heated, behaved as metal to 70° or 80°, that is to say, they lost in conducting power up to that temperature, where it then began to increase. The temperature of this turning-point became lower after each day's heating, until, as in bars I. and II., it is below the lowest temperature at which observations were made.

The behaviour, therefore, of tellurium is intermediate between that of the metal and that of the metalloid; for, according to Hittorf*, selenium increases rapidly in conducting power with the temperature. Graphite and gas-coke † behave in the same manner; and Becquerel; found that gases when heated conduct better than when cold. From these facts we learn another marked difference in the physical properties of the metals and metalloids, viz. that the metals lose in conducting power with an increase of temperature, whereas under the same circumstances the metalloids gain.

In order to be better able to compare the results obtained with the pure metals, we give the following Tables. Table XIV. contains all the formulæ deduced from the observations by the method of least squares, with the conducting power of each metal taken =100 at 0°; Table XV. the mean of the formulæ found for each metal.

^{*} Poggendorff's 'Annalen,' vol. lxxxvi. p. 214.

[†] Philosophical Transactions, 1858, p. 386.

[‡] Ann. de Chim. et de Phys. (iii.) vol. xxxix. p. 388.

TABLE XIV.

	1	1	
	I.	Hard drawn	$\lambda = 100 - 0.38394 \ t + 0.0009723 \ t^2$
	-	Annealed	$\lambda = 100 - 0.37725 \ t + 0.0008900 \ t^2$
~	II.	Hard drawn	$\lambda = 100 - 0.38915 \ t + 0.0010472 \ t^2$
Silver		Annealed	$\lambda = 100 - 0.39287 \ t + 0.0010625 \ t^2$
	III.	Hard drawn	$\lambda = 100 - 0.37544 \ t + 0.0009724 \ t^2$
-	111.	Annealed	$\lambda = 100 - 0.37855 \ t + 0.0009647 \ t^2$
}	T.	Hard drawn	$\lambda = 100 - 0.37351 \ t + 0.0007716 \ t^2$
	•	Annealed	$\lambda = 100 - 0.37291 \ t + 0.0007781 \ t^2$
_	II.	Hard drawn	$\lambda = 100 - 0.39173 \ t + 0.0009394 \ t^2$
Copper	11.	Annealed	$\lambda = 100 - 0.38797 \ t + 0.0008986 \ t^2$
	III.	Hard drawn	$\lambda = 100 - 0.39639 \ t + 0.0010173 \ t^2$
	*1:14	Annealed	$\lambda = 100 - 0.39954 \ t + 0.0010003 \ t^2$
	I.	Hard drawn	$\lambda = 100 - 0.36405 \ t + 0.0008181 \ t^2$
	1.	Annealed	$\lambda = 100 - 0.37017 \ t + 0.0008526 \ t^2$
Gold	II.	Hard drawn	$\lambda = 100 - 0.35877 \ t + 0.0007807 \ t^2$
00.00	III.	Hard drawn	$\lambda = 100 - 0.37551 \ t + 0.0009309 \ t^2$
	111.	Annealed	$\lambda = 100 - 0.36877 \ t + 0.0008390 \ t^2$
اح	I.	Amieaieu	$\lambda = 100 - 0.36845 \ t + 0.0008147 \ t^2$
Zinc	II.	••••••	$\lambda = 100 - 0.37911 \ t + 0.0009058 \ t^2$
Zine	III.	•••••	$\lambda = 100 - 0.36385 \ t + 0.0007618 \ t^2$
<u> </u>	I.		$\lambda = 100 - 0.36745 \ t + 0.0007220 \ t^2$
Cadmium	II.	• • • • • • • • • • • • • • • • • • • •	$\lambda = 100 - 0.39745 t + 0.0007220 t$ $\lambda = 100 - 0.39915 t + 0.0010511 t^2$
Caumium	III.	••••••	$\lambda = 100 - 0.33913$ $t + 0.0010311$ t $\lambda = 100 - 0.33953$ $t + 0.0004995$ t^2
<u> </u>	Ī.		$\lambda = 100 - 0.35867 \ t + 0.0006205 \ t^2$
Tin	II.	••••••	$\lambda = 100 - 0.36730 \ t + 0.0006709 \ t^2$
1 1111	III.		$\lambda = 100 - 0.35492 \ t + 0.0005493 \ t^2$
إ			$\lambda = 100 - 0.38561 \ t + 0.0008989 \ t^2$
Lead	TT		$\lambda = 100 - 0.38553 \ t + 0.0008788 \ t^2$
Leau	III.		$\lambda = 100 - 0.38333 t + 0.0008788 t$ $\lambda = 100 - 0.39153 t + 0.0009661 t^2$
}	I.	•••••••	$\lambda = 100 - 0.33133 t + 0.003301 t$ $\lambda = 100 - 0.41141 t + 0.0011000 t^2$
Arsenic {	II.		$\lambda = 100 - 0.36851 \ t + 0.0006757 \ t^2$
Ĺ	I.		$\lambda = 100 - 0.30851 \ t + 0.000737 \ t$ $\lambda = 100 - 0.39855 \ t + 0.0010332 \ t^2$
Antimony	TT	••••••	$\lambda = 100 - 0.39833 \ t + 0.0010332 \ t$ $\lambda = 100 - 0.37524 \ t + 0.0008477 \ t^2$
21111111011y	III.		$\lambda = 100 - 0.37324 \ t + 0.0008477 \ t^{-1}$ $\lambda = 100 - 0.42099 \ t + 0.0012283 \ t^{2}$
\subseteq	Ι.		$\lambda = 100 - 0.42099 \ t + 0.0012283 \ t$ $\lambda = 100 - 0.36383 \ t + 0.0007293 \ t^2$
Bismuth	II.		$\lambda = 100 - 0.30383 t + 0.0007293 t$ $\lambda = 100 - 0.34146 t + 0.0004593 t^2$
Lisming	III.	••••••	$\lambda = 100 - 0.34140 \ t + 0.0004393 \ t^{-1}$ $\lambda = 100 - 0.35120 \ t + 0.0005298 \ t^{2}$
Ļ	Ι.	••••••	$\lambda = 100 - 0.33120 t + 0.0003298 t$ $\lambda = 100 - 0.074047t + 0.00008672t^2$
Mercury {	II.	•••••••	$\lambda = 100 - 0.074047t + 0.00008072t$ $\lambda = 100 - 0.074820t + 0.00007844t^2$
L.	11.	••••••	V=100-0.0140500±0.000019440
·			

TABLE XV.

Silver Copper Gold Zinc Cadmium Tin Lead Arsenic Antimony Bismuth	$\begin{array}{c} \lambda \! = \! 100 \! - \! 0 \! \cdot \! 38287t \! + \! 0 \! \cdot \! 0009848t^2 \\ \lambda \! = \! 100 \! - \! 0 \! \cdot \! 38701t \! + \! 0 \! \cdot \! 0009009t^2 \\ \lambda \! = \! 100 \! - \! 0 \! \cdot \! 36745t \! + \! 0 \! \cdot \! 0008443t^2 \\ \lambda \! = \! 100 \! - \! 0 \! \cdot \! 37047t \! + \! 0 \! \cdot \! 0008274t^2 \\ \lambda \! = \! 100 \! - \! 0 \! \cdot \! 36871t \! + \! 0 \! \cdot \! 0007575t^2 \\ \lambda \! = \! 100 \! - \! 0 \! \cdot \! 36029t \! + \! 0 \! \cdot \! 0006136t^2 \\ \lambda \! = \! 100 \! - \! 0 \! \cdot \! 38756t \! + \! 0 \! \cdot \! 0009146t^2 \\ \lambda \! = \! 100 \! - \! 0 \! \cdot \! 38996t \! + \! 0 \! \cdot \! 0008879t^2 \\ \lambda \! = \! 100 \! - \! 0 \! \cdot \! 39826t \! + \! 0 \! \cdot \! 0010364t^2 \\ \lambda \! = \! 100 \! - \! 0 \! \cdot \! 35216t \! + \! 0 \! \cdot \! 0005728t^2 \end{array}$
Mean of the above	$\lambda = 100 - 0.37647t + 0.0008340t^{2}$

From the last Table we see how closely the values found for the constants y or z agree together; and to show this more clearly, the conducting powers calculated from these formulæ for 0°, 20°, 40°, 60°, 80°, and 100° are given in Table XVI., together with values calculated from the mean of all the formulæ.

\mathbf{T}	ART	T	X	77	Γ
	ABL	ı Ni		v	

T.	Silver.	Copper.	Gold.	Zinc.	Cadmium.	Tin.	Lead.	Arsenic.	Antimony.	Bismuth.	Calculated values from mean of formulæ.	differ-
0	100·00	100·00	100·00	100·00	100·00	100·00	100·00	100.00	100·00	100·00	100·00	0.00
20	92·74	92·62	92·99	92·92	92·93	93·04	92·62	92.56	92·45	93·18	92·80	0.38
40	86·26	85·96	86·65	86·50	86·46	86·51	85·96	85.82	85·73	86·83	86·27	0.56
60	80·57	80·01	80·98	80·75	80·60	80·59	80·04	79.80	79·84	80·93	80·41	0.61
80	75·67	74·80	76·01	75·66	75·35	75·10	74·85	74.50	74·77	75·49	75·23	0.78
100	71·56	70·31	71·70	71·23	70·70	70·11	70·39	69.88	70·54	70·51	70·69	1.01

Again, in Table XVII., we give the conducting power of the metals compared with hard-drawn silver wire at $0^{\circ}=100$, first at 0° and then at 100° , and, lastly, taking silver at $100^{\circ}=100$.

TABLE XVII.

	Conducti	ng power	Taking silver
	At 0°.	At 100°.	=100 at 100°.
Silver (hard drawn)	100.00	71.56	100.00
Copper (hard drawn)	99.95	70.27	98.20
Gold (hard drawn)	77.96	55.90	78.11
Zinc	29.02	20.67	28.89
Cadmium	23.72	16.77	23.44
Tin	12.36	8.67	12.12
Lead	8.32	5.86	8.18
Arsenic	4.76	3.33	4.65
Antimony	4.62	3.26	4.55
Bismuth	1.245	0.878	1.227

From these Tables we think we may deduce the law, that all pure metals in a solid state vary in conducting power between 0° and 100° to the same extent, more especially as we find that wires of one and the same metal show almost the same differences as were found between the mean results obtained for the different metals. In Table XVIII. two examples of this are given.

TABLE XVIII.

	Cor	oper.	Cadmium.			
T.	I. annealed.	III. annealed.	II.	III.		
0 20 40 60 80	100·00 92·85 86·33 80·43 75·15 70·49	100·00 92·41 85·62 79·63 74·44 70·05	100·00 92·44 85·72 79·84 74·79 70·60	100·00 93·41 87·22 81·42 76·03 71·04		

In Table XIX. the resistances of the copper wires 1, 2, and 3, and those calculated from the mean of all the formulæ, are given; we do this to show that the resistance of

a wire does not increase in direct ratio to the temperature (as stated by some experimenters in this direction), but, on the contrary, the formula for correction of the resistance of a wire for temperature is

and not

$$r = x + yt + zt^2,$$

$$r = x + yt.$$

TABLE XIX.

Firs	st wire, hard	drawn.	Fi	rst wire, ann	ealed.	Secon	d wire, hard	l drawn.	Sec	ond wire, an	nealed.
T.	Resistance.	Increase of resistance for 1°.	Ť.	Resistance.	Increase of resistance for 1°.	т.	Resistance.	Increase of resistance for 1°.	т.	Resistance.	Increase of resistance for 1°.
0 16.86 29.88 51.03 69.52 83.77 98.60	98·382 104·74 109·81 118·72 126·59 132·60 139·22	0·3771 0·3825 0·3985 0·4057 0·4085 0·4142	0 17·0 29·63 50·22 69·60 83·42 99·37	98·241 104·67 109·54 118·08 126·23 132·21 139·10	0·3782 0·3813 0·3950 0·4021 0·4072 0·4112	0 19·17 30·95 48·53 69·22 83·77 99·00	98·412 105·98 110·88 118·32 127·16 133·31 139·80	0.3948 0.4028 0.4102 0.4153 0.4166 0.4181	0 18.96 31.86 52.05 70.27 83.81 99.57	97·902 105·28 110·59 119·08 126·85 132·58 139·36	0.3891 0.3982 0.4069 0.4119 0.4138 0.4164
Thir	d wire, hard	l drawn.	Thi	ird wire, anı	nealed.	Resistance calculated from the mean of the six formulæ found for copper.			Resistance calculated from the mean of all the formulæ.		
0 12.65 25.61 39.52 53.92 69.90 84.87 99.92	109.82 115.72 121.85 128.54 134.82	0·3981 0·4075 0·4134 0·4167 0·4171 0·4175 0·4159	0 13·45 26·15 39·35 55·50 69·90 84·67 99·05	97·806 103·14 108·41 113·99 120·95 127·00 133·25 139·32	0·3966 0·4055 0·4113 0·4170 0·4176 0·4186 0·4190	°0 20 40 60 80 100	100 107·97 116·33 124·98 133·69 142·22	0·3985 0·4082 0·4163 0·4211 0·4222	0 20 40 60 80 100	100 107·76 115·91 124·36 132·92 141·46	0·3880 0·3977 0·4060 0·4115 0·4146

The calculations from a formula of four or more terms, as

$$\lambda = x + yt + zt^2 + at^3,$$

agree better with the observed values than that of three. An example of this is shown in Table XX., where the formulæ, deduced from observations made with a hard-drawn wire (of course previously heated to 100° for several days), of three and four terms, with the differences, are given.

т.	Conducting power.			Conducting	
	Observed.	Calculated from formula of three terms.	Difference.	power, calculated from formula of four terms.	Difference.
10°9 30°1 49°5	95·169 88·537 82·610	95·134 88·588 82·627	+0.035 -0.051 -0.017	95·166 88·534 82·605	$+0.003 \\ +0.003 \\ +0.005$
69·0 82·8 97·9	77·320 73·976 70·579	77·297 73·926 70·619	$+0.023 \\ +0.050 \\ -0.040$	77·304 73·966 70·580	$-0.014 \\ +0.010 \\ -0.001$

TABLE XX.

The formula of three terms, deduced from the observations, was

$$\lambda = 99.137 - 0.37675t + 0.0008728t^2$$

and that of four terms

$$\lambda = 99.307 - 0.39301t + 0.0012318t^2 - 0.000002193t^3$$

From the above it will be seen how much better the observed values agree with the formula of four terms. We have, however, contented ourselves with a formula of three terms, as the conducting powers calculated from it agree with those observed to values corresponding to 0°·1 or 0°·2, and as the calculations for a formula of four terms would have increased the labour of the research to a very great extent. But it may be asked how it happens that the formulæ obtained for wires of one and the same metal vary so much, in fact, show differences almost equal to the mean of those deduced for the different metals?

That this is not due to errors of observation we have repeatedly satisfied ourselves: for compare only the formulæ of the hard-drawn (or rather partially annealed) and the annealed wires, and see how well they agree with each other. It appears, however, to be probably due to the molecular arrangement of the wires being different in each case. Take, for instance, the copper wires experimented with: wire 1 increased in conducting power by heating to 100° for several days, almost to the same extent as if it had been annealed, wire 2 partially so, and wire 3 hardly at all; and here it may be mentioned that silver and copper wires become softer and lose their elasticity, whereas gold does not seem to be annealed at all after having been kept at 100° for several days. Again, take cadmium, where we know that the wires become brittle and crystalline at 80°, and we find the formulæ vary more than those of any other metals; and, lastly, look at the results obtained with bismuth and tellurium, and there can be little doubt that the reason why the formulæ of the wires and bars of the same metal do not agree together is that the molecular arrangement is different in each; and that this is the cause of the differences in the formulæ, we may also assume from the fact that, when the wires on being heated do not at all or only to a very slight degree permanently alter in their conducting power, when cooled again, then the formulæ of wires of the same metal agree very closely with each other. Compare, for instance, those of lead, tin, mercury, &c.

The mean of the conducting powers given in the Tables agrees very well with the mean of the former determinations made with wires of metals of different preparation to that of those used for the experiments described in this paper.

The following questions have suggested themselves during the foregoing investigation, the answers to which we reserve for ourselves. It is intended to make them the subjects of short communications, which from time to time will be laid before the Royal Society:—

- 1. Will a hard-drawn wire become partially annealed by age? and, on the other hand, will an annealed wire become partially hard drawn?
- 2. Will bismuth or tellurium return to their original conducting power in time, or by exposure to intense cold?
- 3. Whether by heating tellurium or any of the metals to a higher temperature than 100° we should not arrive at the same result in a much shorter time.
- 4. What are the thermo-electric properties of bismuth, antimony, tellurium, &c. after being kept at 100° for several days? will they not have altered? It is remarkable that bismuth, which stands at one end of the thermo-electric series, should gain in conducting power after heating for some days, and that antimony and tellurium, at the other end of the series, should lose, the one slightly, the other, with a much higher thermo-electric number, to a very great extent.
 - 5. Will tellurium conduct better in a melted state than the solid?
- 6. What law do the alloys follow as regards the influence of temperature on their conducting power?